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ON THE FLOW OF NON-AXISYMMETRIC PERTURBATIONS OF CYLINDERS VIA SURFACE DIFFUSION

JEREMY LECRONE AND GIERI SIMONETT

ABSTRACT. We study the surface diffusion flow acting on a class of general (non-axisymmetric) perturbations of cylinders \mathcal{C}_r in \mathbb{R}^3 . Using tools from parabolic theory on uniformly regular manifolds, and maximal regularity, we establish existence and uniqueness of solutions to surface diffusion flow starting from (spatially-unbounded) surfaces defined over \mathcal{C}_r via scalar height functions which are uniformly bounded away from the central cylindrical axis. Additionally, we show that \mathcal{C}_r is normally stable with respect to 2π -axially-periodic perturbations if the radius $r > 1$, and unstable if $0 < r < 1$. Stability is also shown to hold in settings with axial Neumann boundary conditions.

1. INTRODUCTION

The surface diffusion flow is a geometric evolution law in which the normal velocity of a moving surface equals the Laplace–Beltrami operator of the mean curvature. More precisely, we assume in the following that Γ_0 is a closed embedded surface in \mathbb{R}^3 . Then the surface diffusion flow is governed by the law

$$(1.1) \quad V(t) = \Delta_{\Gamma(t)} H_{\Gamma(t)}, \quad \Gamma(0) = \Gamma_0.$$

Here $\Gamma = \{\Gamma(t) : t \geq 0\}$ is a family of hypersurfaces, $V(t)$ denotes the velocity in the normal direction of Γ at time t , while $\Delta_{\Gamma(t)}$ and $H_{\Gamma(t)}$ stand for the Laplace–Beltrami operator and the mean curvature of $\Gamma(t)$ (the sum of the principal curvatures in our case), respectively. Both the normal velocity and the mean curvature depend on the local choice of the orientation. Here we consider the case where $\Gamma(t)$ is embedded and encloses a region $\Omega(t)$, and we then choose the outward orientation, so that V is positive if $\Omega(t)$ grows and $H_{\Gamma(t)}$ is positive if $\Gamma(t)$ is convex with respect to $\Omega(t)$.

The unknown quantity in (1.1) is the position and the geometry of the surface $\Gamma(t)$, which evolves in time. Hidden in the formulation of the evolution law (1.1) is a nonlinear partial differential equation of fourth order. This will become more apparent below, where the equation is stated more explicitly.

It is an interesting and significant fact that the surface diffusion flow evolves surfaces in such a way that the volume enclosed by $\Gamma(t)$ is preserved, while the surface area decreases (provided, of course, these quantities are finite). This follows from the well-known relationships

$$(1.2) \quad \frac{d}{dt} \text{Vol}(t) = \int_{\Gamma(t)} V(t) d\sigma = \int_{\Gamma(t)} \Delta_{\Gamma(t)} H_{\Gamma(t)} d\sigma = 0$$

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and

$$(1.3) \quad \begin{aligned} \frac{d}{dt}A(t) &= \int_{\Gamma(t)} V(t)H_{\Gamma(t)} d\sigma = \int_{\Gamma(t)} (\Delta_{\Gamma(t)}H_{\Gamma(t)}) H_{\Gamma(t)} d\sigma \\ &= - \int_{\Gamma(t)} \left| \text{grad}_{\Gamma(t)}H_{\Gamma(t)} \right|_{\Gamma(t)}^2 d\sigma \leq 0, \end{aligned}$$

where $\text{Vol}(t)$ denotes the volume of $\Omega(t)$ and $A(t)$ the surface area of $\Gamma(t)$, respectively. Hence, the preferred ultimate states for the surface diffusion flow in the absence of geometric constraints are spheres. It is also interesting to note that the surface diffusion flow can be viewed as the H^{-1} -gradient flow of the area functional, a fact that was first observed in [20]. This particular structure has been exploited in [29, 30] for devising numerical simulations.

The mathematical equations modeling surface diffusion go back to a paper by Mullins [32] from the 1950s, who was in turn motivated by earlier work of Herring [24]. Since then, the surface diffusion flow (1.1) has received wide attention in the mathematical community (and also by scientists in other fields), see for instance Asai [6], Baras, Duchon, and Robert [8], Bernoff, Bertozzi, and Witelski [9], Cahn and Taylor [11], Cahn, Elliott, and Novick-Cohen [10], Davi and Gurtin [14], Elliott and Garcke [17], Escher, Simonett, and Mayer [18], Escher and Mucha [19], Koch and Lamm [25], LeCrone and Simonett [27], McCoy, Wheeler, and Williams [31], and Wheeler [38, 39].

The primary focus of this article is the surface diffusion flow (1.1) starting from initial surfaces Γ_0 which are perturbations of an infinite cylinder

$$\mathcal{C}_r := \{(x, y, z) \in \mathbb{R}^3 : y^2 + z^2 = r^2, x \in \mathbb{R}\}$$

with radius $r > 0$, centered about the x -axis.

The main results of this paper address

- (a) existence, uniqueness, and regularity of solutions for (1.1) for initial surfaces Γ_0 that are (sufficiently smooth) perturbations of \mathcal{C}_r ;
- (b) existence, uniqueness, and regularity of solutions for (1.1) for initial surfaces Γ_0 that are (sufficiently smooth) perturbations of \mathcal{C}_r subject to periodic or Neumann-type boundary conditions in axial direction;
- (c) nonlinear stability and instability results for \mathcal{C}_r with respect to perturbations subject to periodic or Neumann-type boundary conditions in axial direction.

In section 2 we derive an explicit formulation of (1.1) for surfaces which are parameterized over \mathcal{C}_r in normal direction by height functions $\rho : \mathcal{C}_r \rightarrow \mathbb{R}$, see (2.5). Equation (2.5) is a fourth-order, quasilinear, parabolic partial differential equation for which well-posedness in the unbounded setting can be established by virtue of recent results for maximal regularity on uniformly regular Riemannian manifolds (c.f. [36, 2]). A proof of well-posedness is carried out in section 3 of the current article. For other recent results regarding geometric evolution equations in unbounded settings, refer to Giga, Seki, Umeda [21, 22], wherein the mean curvature flow is shown to close open ends of non-compact surfaces of revolution given appropriate decay rates at the ends of the initial surface. The current setting differs from that of Giga et. al. in many ways; most notably, we consider surfaces which lack symmetry about the cylindrical axis; we control behavior at space infinity via $bc^{2,\alpha}$ regularity bounds (so surfaces do not become radially unbounded); and we consider

the dynamics of surfaces which are uniformly bounded away from the cylindrical axis.

In [27], we studied the surface diffusion flow for the special case of *axisymmetric* and 2π -periodic perturbations Γ_0 of \mathcal{C}_r . In this case it was shown that equilibria of (1.1) consist exactly of cylinders, $2k\pi$ -periodic unduloids and nodoids. In addition, we established nonlinear stability and instability results for cylinders as well as bifurcation results, where r serves as a bifurcation parameter. More specifically, it was shown that cylinders \mathcal{C}_r are stable for $r > 1$, unstable for $r < 1$, and that a subcritical bifurcation of equilibria occurs at $r = 1$, with the bifurcating branch consisting exactly of the family of 2π -periodic unduloids.

In this paper we show that the stability and instability results of [27] remain true for non-axisymmetric perturbations subject to periodic or Neumann-type boundary conditions in axial direction. In particular, we show that small 2π -periodic perturbations Γ_0 of a cylinder \mathcal{C}_r with $r > 1$ exist globally and converge exponentially fast to another nearby cylinder

$$C(\bar{y}, \bar{z}, \bar{r}) := \{(x, y, z) \in \mathbb{R}^3 : (y - \bar{y})^2 + (z - \bar{z})^2 = \bar{r}^2, x \in \mathbb{R}\}.$$

It will be shown that the radius \bar{r} is uniquely determined by the volume of the solid enclosed by Γ_0 and bounded by the planes $x = 0$ and $x = 2\pi$. In order to prove this result, we shall show that all equilibria Γ of (1.1) which are 2π -periodic in axial direction and sufficiently close to \mathcal{C}_r (for $r > 1$) are cylinders $C(\bar{y}, \bar{z}, \bar{r})$, with $(\bar{y}, \bar{z}, \bar{r})$ close to $(0, 0, r)$. This result, which is interesting by itself, is obtained by a center-manifold argument. In order to prove the stability result, we demonstrate that every cylinder \mathcal{C}_r with radius $r > 1$ is *normally stable* with respect to 2π -periodic, $bc^{2+\alpha}$ -regular perturbations. As discussed in section 5, these stability results also encompass perturbations satisfying Neumann-type boundary conditions in x -direction.

2. SURFACE DIFFUSION FLOW NEAR CYLINDERS

Throughout this section, take $r > 0$ fixed, and let $\mathcal{C}_r \subset \mathbb{R}^3$ be the (unbounded) cylinder of radius r which is symmetric about the x -axis. Specifically,

$$\mathcal{C}_r := \{(x, r \cos(\theta), r \sin(\theta)) : x \in \mathbb{R}, \theta \in \mathbb{T}\},$$

where $\mathbb{T} := [0, 2\pi]$ denotes the one-dimensional torus with 0 and 2π identified and \mathbb{T} is equipped with the *periodic* topology generated by the metric

$$d_{\mathbb{T}}(x, y) := \min\{|x - y|, 2\pi - |x - y|\}, \quad x, y \in \mathbb{T}.$$

Further, we equip \mathcal{C}_r with the Riemannian metric $g := dx^2 + r^2 d\theta^2$, inherited from \mathbb{R}^3 via embedding.

Remarks 2.1. Two properties of \mathcal{C}_r are easy to see immediately:

(a) \mathcal{C}_r is a *uniformly regular Riemannian manifold*, as defined by Amann [4]. Indeed, this follows from [4, formula (3.3), Corollary 4.3, and Theorem 3.1], noting that \mathcal{C}_r is realized as the Cartesian product of the circle rS^1 and \mathbb{R} . On a side line we note that the class of uniformly regular (closed) Riemannian manifolds coincides with the class of closed, complete manifolds of bounded geometry, see [15].

(b) \mathcal{C}_r is an equilibrium of (1.1) for every radius $r > 0$.

Indeed, noting that the mean curvature $\mathcal{H}_{\mathcal{C}_r} \equiv 1/r$ is constant throughout the manifold, it follows that

$$\Delta_{\mathcal{C}_r} \mathcal{H}_{\mathcal{C}_r} \equiv 0,$$

so that the normal velocity $V \equiv 0$ and the surface remains fixed in space.

Now, consider a scalar-valued function $\rho : \mathcal{C}_r \rightarrow \mathbb{R}$ such that

$$\rho(p) > -r, \quad \text{for all } p \in \mathcal{C}_r,$$

and define a new surface $\Gamma = \Gamma(\rho)$ over the reference manifold \mathcal{C}_r as

$$\Gamma := \{p + \rho(p)\nu(p) : p \in \mathcal{C}_r\},$$

where ν denotes the unit outer normal field over \mathcal{C}_r . We also note that the points $p \in \mathcal{C}_r$ are uniquely determined via the surface coordinates $(x, \theta) \in \mathbb{R} \times \mathbb{T}$, and so we also view the height function ρ as a function of the variables x and θ . Thus, abusing notation, we will interchangeably refer to $\rho(p)$ and $\rho(x, \theta)$, where, in fact,

$$\rho(x, \theta) := \rho((x, r \cos(\theta), r \sin(\theta))), \quad (x, \theta) \in \mathbb{R} \times \mathbb{T}.$$

Extending this slight abuse of notation, one easily identifies an explicit parametrization for $\Gamma(\rho)$ via

$$(2.1) \quad \varphi(x, \theta) := (x, r \cos(\theta), r \sin(\theta)) + \rho(x, \theta)(0, \cos(\theta), \sin(\theta)), \quad (x, \theta) \in \mathbb{R} \times \mathbb{T}.$$

With the above construction, we have a correspondence between height functions $\rho > -r$ and a particular class of embedded manifolds $\Gamma(\rho) \subset \mathbb{R}^3$ which we refer to as *axially-definable* — namely, those manifolds which admit an axial parametrization $\varphi > 0$ of the form (2.1). Within this class of manifolds, the geometry, time-dependent evolution, and regularity are determined by the height function ρ itself. We turn to the task of explicitly expressing these relationships, and deriving the governing equation for surface diffusion, in the following subsection. For the purpose of expressing the necessary geometric quantities, we assume for the moment that $\rho = \rho(x, \theta)$ is smooth enough for all following calculations to work. The precise desired regularity of ρ will be addressed in detail when we discuss existence and uniqueness of solutions.

2.1. Geometry of $\Gamma(\rho)$. Utilizing explicit parametrizations of the form (2.1), we find the coefficients of the first fundamental form of $\Gamma(\rho)$ as

$$g_{11} = (1 + \rho_x^2), \quad g_{12} = g_{21} = \rho_x \rho_\theta, \quad g_{22} = ((r + \rho)^2 + \rho_\theta^2),$$

and the metric is thus given by $g = g_{11} dx^2 + 2g_{12} dx d\theta + g_{22} d\theta^2$. Then,

$$g^{11} = \frac{(r + \rho)^2 + \rho_\theta^2}{\mathcal{G}}, \quad g^{12} = g^{21} = \frac{-\rho_x \rho_\theta}{\mathcal{G}}, \quad g^{22} = \frac{1 + \rho_x^2}{\mathcal{G}},$$

where the matrix $[g^{ij}]$ is the inverse of $[g_{ij}]$ and

$$(2.2) \quad \mathcal{G} = \mathcal{G}(\rho) := \det(g_{ij}) = (r + \rho)^2(1 + \rho_x^2) + \rho_\theta^2.$$

Likewise, the second fundamental form is expressed via the coefficients

$$\begin{aligned}\mathbb{I}_{11} &= \frac{(r + \rho)\rho_{xx}}{\sqrt{\mathcal{G}}}, \\ \mathbb{I}_{12} = \mathbb{I}_{21} &= \frac{(r + \rho)\rho_{x\theta} - \rho_x\rho_\theta}{\sqrt{\mathcal{G}}}, \\ \mathbb{I}_{22} &= \frac{(r + \rho)(\rho_{\theta\theta} - (r + \rho)) - 2\rho_\theta^2}{\sqrt{\mathcal{G}}}.\end{aligned}$$

The mean curvature $\mathcal{H}(\rho) := \mathcal{H}_{\Gamma(\rho)}$ and Laplace–Beltrami operator $\Delta_{\Gamma(\rho)}$ are expressed using well-known formulas, c.f. [16, 34]. From [34, Equation (5)],

$$(2.3) \quad \mathcal{H}(\rho) = \frac{\rho_\theta^2 - (r + \rho)[(r + \rho)^2 + \rho_\theta^2]\rho_{xx} + (1 + \rho_x^2)\rho_{\theta\theta} - 2\rho_x\rho_\theta\rho_{x\theta}}{\mathcal{G}^{3/2}} + \frac{1}{\mathcal{G}^{1/2}},$$

and, from [34, Section 2.7],

$$\Delta_{\Gamma(\rho)} = \frac{1}{\sqrt{\mathcal{G}}} \left\{ \partial_x \left[\frac{(r + \rho)^2 + \rho_\theta^2}{\sqrt{\mathcal{G}}} \partial_x - \frac{\rho_x\rho_\theta}{\sqrt{\mathcal{G}}} \partial_\theta \right] + \partial_\theta \left[\frac{1 + \rho_x^2}{\sqrt{\mathcal{G}}} \partial_\theta - \frac{\rho_x\rho_\theta}{\sqrt{\mathcal{G}}} \partial_x \right] \right\}.$$

The normal velocity of the surface $\Gamma(\rho) = \Gamma(\rho(t))$ is likewise

$$V_{\Gamma(\rho)} = \frac{(r + \rho)\rho_t}{\sqrt{\mathcal{G}(\rho)}}.$$

Throughout, we often streamline notation by omitting explicit reference to dependence upon the spatial or temporal variables; e.g. $\rho = \rho(t) = \rho(x, \theta) = \rho(t, p) = \rho(t, x, \theta)$ and $\Gamma(\rho) = \Gamma(\rho(t))$, etc.

2.2. The Surface Diffusion Flow. With the formulas from the last subsection, we can now express the governing equation for surface diffusion of $\Gamma(\rho)$ as an evolution equation for the height function ρ alone. Thus, defining the (formal) operator

$$(2.4) \quad \begin{aligned}G(\rho) := \frac{1}{(r + \rho)} & \left\{ \partial_x \left[\frac{(r + \rho)^2 + \rho_\theta^2}{\sqrt{\mathcal{G}(\rho)}} \partial_x \mathcal{H}(\rho) - \frac{\rho_x\rho_\theta}{\sqrt{\mathcal{G}(\rho)}} \partial_\theta \mathcal{H}(\rho) \right] \right. \\ & \left. + \partial_\theta \left[\frac{1 + \rho_x^2}{\sqrt{\mathcal{G}(\rho)}} \partial_\theta \mathcal{H}(\rho) - \frac{\rho_x\rho_\theta}{\sqrt{\mathcal{G}(\rho)}} \partial_x \mathcal{H}(\rho) \right] \right\},\end{aligned}$$

we arrive at the expression

$$(2.5) \quad \begin{cases} \rho_t(t, p) = [G(\rho(t))](p), & \text{for } t > 0, p \in \mathcal{C}_r, \\ \rho(0) = \rho_0, & \text{on } \mathcal{C}_r, \end{cases}$$

for the surface diffusion flow for axially-definable surfaces $\Gamma(\rho)$.

Remark 2.2. The notation $[G(\rho(t))](p)$ reflects the functional-analytic framework we use to address equation (2.5). For each $t \geq 0$, we consider $\rho(t)$ as an element of an appropriate Banach space of regular functions defined on \mathcal{C}_r . Then, G maps $\rho(t)$ to another function defined over \mathcal{C}_r which we then evaluate pointwise using the notation $[G(\rho(t))](p)$.

3. EXISTENCE AND UNIQUENESS OF SOLUTIONS

In this section we establish existence and uniqueness of solutions for (2.5). One essential tool that we use throughout is the property of maximal regularity, also called *optimal regularity*. Maximal regularity has received considerable attention in connection with nonlinear parabolic partial differential equations, c.f. [1, 5, 12, 28, 33].

3.1. Maximal Regularity. Although maximal regularity can be developed in a more general setting, we will focus on the setting of *continuous* maximal regularity and direct the interested reader to the references [1, 28] for a general development of the theory.

Let $\mu \in (0, 1]$, $J := [0, T]$, for some $T > 0$, and let E be a (real or complex) Banach space. Following the notation of [12], we define spaces of continuous functions on $\dot{J} := J \setminus \{0\}$ with prescribed singularity at 0 as

$$(3.1) \quad \begin{aligned} BUC_{1-\mu}(J, E) &:= \left\{ u \in C(\dot{J}, E) : [t \mapsto t^{1-\mu}u(t)] \in BUC(\dot{J}, E) \text{ and} \right. \\ &\quad \left. \lim_{t \rightarrow 0^+} t^{1-\mu}\|u(t)\|_E = 0 \right\}, \quad \mu \in (0, 1) \\ \|u\|_{B_{1-\mu}} &:= \sup_{t \in J} t^{1-\mu}\|u(t)\|_E, \end{aligned}$$

where BUC denotes the space consisting of bounded, uniformly continuous functions. We also define the subspace

$$BUC_{1-\mu}^1(J, E) := \left\{ u \in C^1(\dot{J}, E) : u, \dot{u} \in BUC_{1-\mu}(J, E) \right\}, \quad \mu \in (0, 1)$$

and we set

$$BUC_0(J, E) := BUC(J, E), \quad BUC_0^1(J, E) := BUC^1(J, E).$$

If $J = [0, a)$ for $a > 0$, then we set

$$\begin{aligned} C_{1-\mu}(J, E) &:= \{u \in C(\dot{J}, E) : u \in BUC_{1-\mu}([0, T], E), \quad T < \sup J\}, \\ C_{1-\mu}^1(J, E) &:= \{u \in C^1(\dot{J}, E) : u, \dot{u} \in C_{1-\mu}(J, E)\}, \quad \mu \in (0, 1], \end{aligned}$$

which we equip with the natural Fréchet topologies induced by $BUC_{1-\mu}([0, T], E)$ and $BUC_{1-\mu}^1([0, T], E)$, respectively.

If E_1 and E_0 are a pair of Banach spaces such that E_1 is continuously embedded in E_0 , denoted $E_1 \hookrightarrow E_0$, we set

$$(3.2) \quad \begin{aligned} \mathbb{E}_0(J) &:= BUC_{1-\mu}(J, E_0), \quad \mu \in (0, 1], \\ \mathbb{E}_1(J) &:= BUC_{1-\mu}^1(J, E_0) \cap BUC_{1-\mu}(J, E_1), \end{aligned}$$

where $\mathbb{E}_1(J)$ is a Banach space with the norm

$$\|u\|_{\mathbb{E}_1(J)} := \sup_{t \in \dot{J}} t^{1-\mu} \left(\|\dot{u}(t)\|_{E_0} + \|u(t)\|_{E_1} \right).$$

It follows that the trace operator $\gamma : \mathbb{E}_1(J) \rightarrow E_0$, defined by $\gamma v := v(0)$, is well-defined and we denote by $\gamma \mathbb{E}_1$ the image of γ in E_0 , which is itself a Banach space when equipped with the norm

$$\|x\|_{\gamma \mathbb{E}_1} := \inf \left\{ \|v\|_{\mathbb{E}_1(J)} : v \in \mathbb{E}_1(J) \text{ and } \gamma v = x \right\}.$$

Given $B \in \mathcal{L}(E_1, E_0)$, closed as an operator on E_0 , we say $(\mathbb{E}_0(J), \mathbb{E}_1(J))$ is a pair of maximal regularity for B and write $B \in \mathcal{MR}_\mu(E_1, E_0)$, if

$$\left(\frac{d}{dt} + B, \gamma \right) \in \mathcal{L}_{\text{isom}}(\mathbb{E}_1(J), \mathbb{E}_0(J) \times \gamma\mathbb{E}_1), \quad \mu \in (0, 1),$$

where $\mathcal{L}_{\text{isom}}$ denotes the set of bounded linear isomorphisms. In particular, $B \in \mathcal{MR}_\mu(E_1, E_0)$ if and only if for every $(f, u_0) \in \mathbb{E}_0(J) \times \gamma\mathbb{E}_1$, there exists a unique solution $u \in \mathbb{E}_1(J)$ to the inhomogeneous Cauchy problem

$$\begin{cases} \dot{u}(t) + Bu(t) = f(t), & t \in J, \\ u(0) = u_0. \end{cases}$$

Moreover, in the current setting, it follows that $\gamma\mathbb{E}_1 \doteq (E_0, E_1)_{\mu, \infty}^0$, i.e. the trace space $\gamma\mathbb{E}_1$ is topologically equivalent to the noted continuous interpolation spaces of Da Prato and Grisvard, c.f. [1, 12, 13, 28].

We shall establish well-posedness of (2.5) in the setting of *little-Hölder* regular height functions ρ . These results seem to constitute the first existence and uniqueness results for the surface diffusion flow acting on unbounded (closed) surfaces. Note that we enforce minimal conditions on the ends of the initial surface, namely we look at surfaces which are uniformly bounded away from the axis of definition (i.e. the x -axis in our current setting) and satisfy minimal regularity assumptions.

For the convenience of the reader, we include a brief definition of the little-Hölder spaces on \mathcal{C}_r . For $\alpha \in (0, 1)$ and $U \subset \mathbb{R}^n$ open, we define $bc^\alpha(U)$ to be the closure of the bounded smooth functions $BC^\infty(U)$ in the topology of the bounded Hölder functions $BC^\alpha(U)$. Further, for $k \in \mathbb{N}$, we define $bc^{k+\alpha}(U)$ to be the space of k -times continuously differentiable functions such that the k^{th} -order derivatives are in $bc^\alpha(U)$. We then define the space of $(k + \alpha)$ little-Hölder regular functions on \mathcal{C}_r , $bc^{k+\alpha}(\mathcal{C}_r)$, via an atlas of local charts and a subordinate localization system. For more details regarding function spaces on uniformly regular (and singular) Riemannian manifolds see [2, 36].

3.2. Quasilinear Structure and Maximal Regularity. Expanding terms of (2.4), it is straight-forward to see that $G(\rho)$ is a fourth-order quasilinear operator of the form

$$G(\rho) = -\mathcal{A}(\rho)\rho + F(\rho) := - \left(\sum_{|\beta|=3,4} b_\beta(\rho, \partial^1 \rho, \partial^2 \rho) \partial^\beta \rho \right) + F(\rho, \partial^1 \rho, \partial^2 \rho),$$

where $\beta = (\beta_1, \beta_2) \in \mathbb{N}^2$ is a multi-index, $|\beta| := \beta_1 + \beta_2$ its length, $\partial^\beta = \partial^{(\beta_1, \beta_2)} := \partial_x^{\beta_1} \partial_\theta^{\beta_2}$, and $\partial^k \rho$, for $k \in \mathbb{N}$, denotes the collection of all derivatives $\partial^\beta \rho$ for $|\beta| = k$. This quasilinear structure plays an important role in our well-posedness results below, for which we must look more closely at the fine properties of the variable coefficient linear operator $\mathcal{A}(\rho)$.

Expanding terms, we have the following highest-order variable coefficients for the linear operator $\mathcal{A}(\rho)$:

$$\begin{aligned} b_{(4,0)}(\rho) &= \frac{((r+\rho)^2 + \rho_\theta^2)^2}{\mathcal{G}^2}, \\ b_{(3,1)}(\rho) &= -\frac{4\rho_x\rho_\theta((r+\rho)^2 + \rho_\theta^2)}{\mathcal{G}^2}, \\ b_{(2,2)}(\rho) &= \frac{2((r+\rho)^2 + \rho_\theta^2)(1 + \rho_x^2) + 4\rho_x^2\rho_\theta^2}{\mathcal{G}^2}, \\ b_{(1,3)}(\rho) &= -\frac{4\rho_x\rho_\theta(1 + \rho_x^2)}{\mathcal{G}^2}, \\ b_{(0,4)}(\rho) &= \frac{(1 + \rho_x^2)^2}{\mathcal{G}^2}. \end{aligned}$$

Therefore, the principal symbol $\sigma[\mathcal{A}(\rho)]$ of the linear operator \mathcal{A} satisfies

$$\begin{aligned} \sigma[\mathcal{A}(\rho)](p, \xi) &:= \sum_{j+k=4} [b_{(j,k)}(\rho)](p)(i\xi_1)^j(i\xi_2)^k \\ &= \sum_{j+k=4} [b_{(j,k)}(\rho)](p)\xi_1^j\xi_2^k \\ (3.3) \quad &= \frac{1}{\mathcal{G}^2} \left(((r+\rho)^2 + \rho_\theta^2)\xi_1^2 + (1 + \rho_x^2)\xi_2^2 - 2\rho_x\rho_\theta\xi_1\xi_2 \right)^2 \\ &\geq \frac{1}{\mathcal{G}^2} \left((r+\rho)^2\xi_1^2 + \xi_2^2 \right)^2. \end{aligned}$$

With these quantities explicitly expressed, we will show that \mathcal{A} satisfies the property of continuous maximal regularity on \mathcal{C}_r . Well-posedness of (2.5) then follows by exploiting the quasilinear structure of the parabolic equation.

Proposition 3.1 (maximal regularity). *Fix $\varepsilon > 0$, $\alpha \in (0, 1)$, and take $\mu \in [1/2, 1]$ so that $4\mu + \alpha \notin \mathbb{Z}$. Further, let $V_\mu := bc^{4\mu+\alpha}(\mathcal{C}_r) \cap [\rho > \varepsilon - r]$ be the set of admissible height functions. Then it follows that*

$$(\mathcal{A}, F) \in C^\omega \left(V_\mu, \mathcal{M}_\nu(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r)) \times bc^\alpha(\mathcal{C}_r) \right),$$

where C^ω denotes the space of real analytic mappings between (open subsets of) Banach spaces.

Proof. The regularity of $\mathcal{A}(\cdot)$ and $F(\cdot)$ is a consequence of the analyticity of the following mappings: (which is easy to show using standard methods in nonlinear analysis and theory of function spaces)

$$\begin{aligned} [\rho \mapsto 1/\rho] &: V_0 \rightarrow bc^\alpha(\mathcal{C}_r), \\ [\rho \mapsto \partial^\beta \rho] &: bc^{\sigma+|\beta|}(\mathcal{C}_r) \rightarrow bc^\sigma(\mathcal{C}_r), \\ [(\rho, h) \mapsto \rho h] &: bc^\alpha(\mathcal{C}_r) \times bc^\alpha(\mathcal{C}_r) \rightarrow bc^\alpha(\mathcal{C}_r), \end{aligned}$$

where $V_0 := bc^\alpha(\mathcal{C}_r) \cap [\rho > \varepsilon - r]$. Additionally, one notes that the regularity of $\mathcal{A} : V_\mu \rightarrow \mathcal{L}(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r))$ is determined by the regularity of the coefficients $b_\beta : V_\mu \rightarrow bc^\alpha(\mathcal{C}_r)$. Then, one is left only to show that $\mathcal{A}(\rho)$ is in the denoted maximal regularity class \mathcal{M}_ν for $\rho \in V_\mu$. Note that $\mathcal{M}_\nu(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r))$ is an open subset of $\mathcal{L}(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r))$ (by [12, Lemma 2.5(a)], for instance); thus the

regularity of \mathcal{A} mapping into $\mathcal{M}_\nu(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r))$ is determined by the regularity of \mathcal{A} mapping into $\mathcal{L}(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r))$.

To show $\mathcal{A}(\rho) \in \mathcal{M}_\nu(bc^{4+\alpha}(\mathcal{C}_r), bc^\alpha(\mathcal{C}_r))$, by [36, Theorem 3.6], it suffices to show that $\mathcal{A}(\rho)$ is uniformly strongly elliptic on \mathcal{C}_r ; i.e. we need to show that there exist constants $0 < c_1 < c_2$ such that

$$c_1 \leq \operatorname{Re}(\sigma[\mathcal{A}(\rho)](p, \xi)) \leq c_2, \quad \text{for } (p, \xi) \in \mathcal{C}_r \times \mathbb{R}^2, \quad \text{with } |\xi| = 1.$$

However, these bounds are obvious from the expression (3.3) and the assumption that admissible functions ρ are uniformly bounded from below and above on \mathcal{C}_r . \square

3.3. Well-Posedness of (2.5). We now prove existence and uniqueness of solutions to (2.5). Due to the non-compact setting, we must take care of the behavior of the height function as the axial variable approaches $\pm\infty$. For this purpose, we concentrate on surfaces which remain uniformly bounded away from the x-axis (i.e. height function $\rho > -r + \varepsilon$ for $\rho : \mathcal{C}_r \rightarrow \mathbb{R}$ and $\varepsilon > 0$). Also note that functions in the spaces $bc^\sigma(\mathcal{C}_r)$ are bounded from above by definition.

Proposition 3.2 (existence and uniqueness). *Fix $\varepsilon > 0$, $\alpha \in (0, 1)$, and take $\mu \in [1/2, 1]$ so that $4\mu + \alpha \notin \mathbb{Z}$. For each initial value*

$$\rho_0 \in V_\mu := bc^{4\mu+\alpha}(\mathcal{C}_r) \cap [\rho > \varepsilon - r],$$

there exists a unique maximal solution to (2.5)

$$\rho(\cdot, \rho_0) \in C_{1-\mu}^1(J(\rho_0), bc^\alpha(\mathcal{C}_r)) \cap C_{1-\mu}(J(\rho_0), bc^{4+\alpha}(\mathcal{C}_r)),$$

where $J(\rho_0) := [0, t^+(\rho_0)) \subset \mathbb{R}_+$ denotes the maximal interval of existence for initial data ρ_0 . Further, it follows that

$$\mathcal{D} := \bigcup_{\rho_0 \in V_\mu} J(\rho_0) \times \{\rho_0\}$$

is open in $\mathbb{R}_+ \times V_\mu$ and the map $[(t, \rho_0) \mapsto \rho(t, \rho_0)]$ defines an analytic semiflow on V_μ . Moreover, if the solution $\rho(\cdot, \rho_0)$ satisfies:

- (i) $\rho(\cdot, \rho_0) \in UC(J(\rho_0), bc^{4\mu+\alpha}(\mathcal{C}_r))$, and
- (ii) *there exists $\delta > 0$ so that $\operatorname{dist}_{bc^{4\mu+\alpha}(\mathcal{C}_r)}(\rho(t, \rho_0), \partial V_\mu) > \delta$ for all $t \in J(\rho_0)$.*

Then it must hold that $t^+(\rho_0) = \infty$ and so $\rho(\cdot, \rho_0)$ is a global solution of (2.5).

Remark 3.3. Condition (ii) of the proposition can also be interpreted by explicitly identifying the boundary of the admissible set V_μ . In particular, $\rho(\cdot, \rho_0)$ remains bounded away from ∂V_μ if and only if there exists some $0 < M < \infty$ such that, for all $t \in J(\rho_0)$

- (ii.a) $\rho(t, \rho_0)(p) \geq 1/M + \varepsilon - r$ for all $p \in \mathcal{C}_r$, and
- (ii.b) $\|\rho(t, \rho_0)\|_{bc^{4\mu+\alpha}(\mathcal{C}_r)} \leq M$.

Proof. Existence of maximal solutions is proved in the same way as Proposition 2.2 of [27], whereas the claim for global solutions differs slightly from Proposition 2.3 of [27] due to the fact that we do not have compact embedding of little-Hölder spaces over the non-compact manifold \mathcal{C}_r . In particular, well-posedness follows from [12, Theorems 3.1 and 4.1(c)], [36, Proposition 2.2(c)], and Proposition 3.1, while the semiflow properties follow from [12] when $\mu < 1$ and from [5] in case $\mu = 1$. \square

3.4. Well-Posedness: Axially Periodic Surfaces. In order to address the stability of cylinders under the flow of (2.5) — for which we will use explicit information regarding the spectrum of the linearization of G — we restrict our setting to one within which spectral calculations are tractable.

For $a > 0$ fixed, define the axial shift operator

$$T_a := [x \mapsto x + a] : \mathbb{R} \rightarrow \mathbb{R},$$

which naturally acts on functions $\rho \in bc^\sigma(\mathcal{C}_r)$ as

$$[T_a \rho](x, \theta) = \rho(x + a, \theta).$$

We define

$$bc_{a,per}^\sigma(\mathcal{C}_r) := \{\rho \in bc^\sigma(\mathcal{C}_r) : T_a \rho = \rho\},$$

which we refer to as axially a -periodic little-Hölder functions on \mathcal{C}_r . It is a straightforward exercise to see that $bc_{a,per}^\sigma(\mathcal{C}_r)$ forms a closed subspace of $bc^\sigma(\mathcal{C}_r)$, and is hence a Banach space. We thus consider the properties of the surface diffusion evolution operator G as it acts on the axially periodic spaces.

Proposition 3.4 (G preserves periodicity). *If $\rho \in bc_{a,per}^{4+\alpha}(\mathcal{C}_r) \cap [\rho > -r]$, then $G(\rho) \in bc_{a,per}^\alpha(\mathcal{C}_r)$.*

Proof. It suffices to show that T_a commutes with the nonlinear operator G , since ρ a -periodic would then imply $T_a G(\rho) = G(T_a \rho) \equiv G(\rho)$, as desired. The fact that T_a indeed commutes with G follows directly from the commutativity of T_a with the operations:

$$[\rho \mapsto 1/\rho], \quad [\rho \mapsto \partial^\beta \rho], \quad \text{and} \quad [(\rho, h) \mapsto \rho h],$$

where, as before, $\partial^\beta = \partial^{(\beta_1, \beta_2)} = \partial_x^{\beta_1} \partial_\theta^{\beta_2}$. \square

Remark 3.5. The fact that G preserves periodicity can also be seen directly from the geometric setting. In particular, recall that G was constructed by modeling the evolution of the surface $\Gamma(\rho)$ via (1.1), which depends only upon the geometry of the surface $\Gamma(\rho)$ itself. Thus, if ρ is periodic, then $\Gamma(\rho)$, and all relevant geometric quantities on $\Gamma(\rho)$ must also be periodic, and hence the evolution cannot break this periodicity.

The well-posedness results for (2.5) simplify slightly when restricted to periodic surfaces, owing to the fact that the evolution is determined by the restriction to one interval of periodicity. In essence, periodic surfaces are expressed entirely in a compact setting.

Proposition 3.6 (periodic well-posedness). *Fix $\alpha \in (0, 1)$, $a > 0$ and take $\mu \in [1/2, 1]$ so that $4\mu + \alpha \notin \mathbb{Z}$. For each initial value*

$$\rho_0 \in V_{\mu,per} := bc_{a,per}^{4\mu+\alpha}(\mathcal{C}_r) \cap [\rho > -r],$$

there exists a unique maximal a -periodic solution to (2.5)

$$\rho(\cdot, \rho_0) \in C_{1-\mu}^1(J(\rho_0), bc_{a,per}^\alpha(\mathcal{C}_r)) \cap C_{1-\mu}(J(\rho_0), bc_{a,per}^{4+\alpha}(\mathcal{C}_r)).$$

It follows that

$$\mathcal{D} := \bigcup_{\rho_0 \in V_{\mu,per}} J(\rho_0) \times \{\rho_0\}$$

is open in $\mathbb{R}_+ \times V_{\mu,per}$ and the map $[(t, \rho_0) \mapsto \rho(t, \rho_0)]$ defines an analytic semiflow on $V_{\mu,per}$. Moreover, if there exists $0 < M < \infty$ so that, for all $t \in J(\rho_0)$,

- a) $\rho(t, \rho_0)(p) \geq 1/M - r$ for all $p \in \mathcal{C}_r$, and
 b) $\|\rho(t, \rho_0)\|_{bc_{a,per}^{4\mu+\alpha}(\mathcal{C}_r)} \leq M$,

then it must hold that $t^+(\rho_0) = \infty$ and $\rho(\cdot, \rho_0)$ is a global solution.

Proof. By the assumption $\rho_0 > -r$ and periodicity, it is clear that ρ_0 is uniformly bounded away from the x -axis. According to Theorem 3.2, the surface diffusion flow (2.5) admits a unique solution $\rho(\cdot) = \rho(\cdot, \rho_0)$ on a maximal interval of existence $J(\rho_0)$. It, thus, only remains to show that $\rho(t)$ is a -periodic for each $t \in J(\rho_0)$. Let $\rho_a(t) := T_a\rho(t)$ for $t \in J(\rho_0)$. Then

$$\begin{aligned}\partial_t \rho_a(t) &= T_a \partial_t \rho(t) = T_a G(\rho(t)) = G(T_a \rho(t)), \quad t \in J(\rho_0), \\ \rho_a(0) &= T_a \rho(0) = T_a \rho_0 = \rho_0,\end{aligned}$$

showing that $\rho(\cdot)$ and $\rho_a(\cdot)$ both are solutions of (2.5) with the same initial value. By uniqueness, $T_a \rho(t) = \rho(t)$, implying that $\rho(t)$ is a -periodic for all $t \in J(\rho_0)$.

The global well-posedness result follows from compactness of the embedding $bc_{a,per}^{4+\alpha}(\mathcal{C}_r) \hookrightarrow bc_{a,per}^\alpha(\mathcal{C}_r)$ and [12, Theorem 4.1(d)]. \square

For $\rho \in bc_{a,per}^\alpha(\mathcal{C}_r)$ we define the area of $\Gamma(\rho) = \{p + \rho(p)\nu(p) : p \in \mathcal{C}_r\}$ and the volume of the solid enclosed by $\Gamma(\rho)$, respectively, in a natural way by just considering the part of $\Gamma(\rho)$ on an interval of periodicity in the x -direction. More precisely, we set

$$(3.4) \quad \begin{aligned}\mathcal{C}_{r,a} &= \{(x, y, z) \in \mathbb{R}^3 : y^2 + z^2 = r^2, x \in [0, a]\}, \\ \Gamma(\rho, a) &= \{p + \rho(p)\nu(p) : p \in \mathcal{C}_{r,a}\}.\end{aligned}$$

Then we have the following result.

Proposition 3.7. *Fix $\alpha \in (0, 1)$, $a > 0$ and take $\mu \in [1/2, 1]$ so that $4\mu + \alpha \notin \mathbb{Z}$. For each initial value*

$$\rho_0 \in V_{\mu,per} := bc_{a,per}^{4\mu+\alpha}(\mathcal{C}_r) \cap [\rho > -r],$$

let $\rho = \rho(t, \rho_0)$ be the solution of (2.5). Then the surface diffusion flow preserves the volume of the region enclosed by $\Gamma(\rho(t), a)$ and bounded by the planes $x = 0$ and $x = a$. Moreover, the flow reduces the surface area of $\Gamma(\rho(t), a)$.

Proof. Although the assertions follow from (1.2)-(1.3) it will be instructive to give an independent proof that only relies on basic computations. A short moment of reflection shows that volume is given by

$$(3.5) \quad \text{Vol}(\rho(t)) = \frac{1}{2} \int_{[0,a] \times [0,2\pi]} (r + \rho(t))^2 d\theta dx.$$

Using the short form $\rho = \rho(t, x, \theta)$, and keeping in mind that $\rho(t, x, \theta)$ is periodic in $(x, \theta) \in [0, a] \times [0, 2\pi]$, one immediately obtains from (2.4)

$$\frac{d}{dt} \text{Vol}(\rho(t)) = \int (r + \rho) \rho_t d\theta dx = 0.$$

Next we note that the surface area of $\Gamma(\rho(t), a)$ is given by

$$\text{A}(\rho(t)) = \int_{[0,a] \times [0,2\pi]} \sqrt{\mathcal{G}(\rho(t))} d\theta dx.$$

It follows from integration by parts (using periodicity) and formulas (2.2)–(2.4)

$$\begin{aligned} \frac{d}{dt}A(\rho(t)) &= \int \left[\frac{(r+\rho)(1+\rho_x^2)}{\sqrt{G}} - \partial_x \left(\frac{(r+\rho)^2 \rho_x}{\sqrt{G}} \right) - \partial_\theta \left(\frac{\rho_\theta}{\sqrt{G}} \right) \right] \rho_t \, d\theta dx \\ &= \int (r+\rho) \mathcal{H} \rho_t \, d\theta dx \\ &= - \int \frac{1}{\sqrt{G}} \left[((r+\rho)^2 + \rho_\theta^2) (\mathcal{H}_x)^2 - 2\rho_x \rho_\theta (\mathcal{H}_x \mathcal{H}_\theta + (1+\rho_x^2) (\mathcal{H}_\theta)^2) \right] d\theta dx \\ &\leq - \int \frac{1}{\sqrt{G}} \left[(r+\rho)^2 (\mathcal{H}_x)^2 + (\mathcal{H}_\theta)^2 \right] d\theta dx. \end{aligned}$$

Hence $A(\rho(t))$ is non-increasing and acts as a Lyapunov functional for the evolution in the a -periodic setting. We also conclude that a -periodic equilibria of (2.5) correspond exactly to (a -periodic) surfaces of constant mean curvature. \square

Before moving on to establish dynamic properties of solutions, we state the following characterization for axial periodic height functions defined on C_r . This characterization allows explicit access to calculations involving the linearization and spectrum of $DG(0)$ in the periodic setting.

Proposition 3.8. *Given $\rho \in V_\mu := bc_{a,per}^{4\mu+\alpha}(C_r)$, it follows that ρ satisfies the Fourier series representation*

$$\rho(x, \theta) = \sum_{m,n \in \mathbb{Z}} \hat{\rho}(m, n) e^{2\pi i m x/a} e^{i n \theta}, \quad (x, \theta) \in \mathbb{R} \times \mathbb{T},$$

where

$$\hat{\rho}(m, n) := \frac{1}{2\pi a} \int_{[0,a] \times \mathbb{T}} \rho(x, \theta) e^{-2\pi i m x/a} e^{-i n \theta} \, d\theta dx.$$

Proof. This follows from the realization of ρ as a function over $\mathbb{R} \times \mathbb{T}$ — as discussed in Section 2 — and classic Fourier analysis results for periodic functions in higher dimensions (c.f. [35]). \square

4. NONLINEAR STABILITY AND INSTABILITY OF CYLINDERS UNDER PERIODIC PERTURBATIONS

For the remainder of the paper we assume $a = 2\pi$, and consider the behavior of the evolution equation (2.5) acting on 2π -axial-periodic admissible height functions defined over C_r (i.e. $\rho_0 \in bc_{2\pi,per}^{2+\alpha}(C_r)$ and $\rho_0 > -r$). Specifically, we are interested in the nonlinear dynamics of (2.5) in a neighborhood of $\rho \equiv 0$ (i.e. near the cylinder C_r itself). Thus, we begin in the next subsection by determining the spectral properties of the linearized evolution operator $DG(0)$.

4.1. Linearization of G . Utilizing the expression (2.4), one readily computes the Fréchet derivative of G , which simplifies considerably upon evaluation at the constant function $\rho \equiv 0$ to the operator $DG(0) \in \mathcal{L}(bc_{2\pi,per}^{4+\alpha}(C_r), bc_{2\pi,per}^\alpha(C_r))$ given by

$$(4.1) \quad DG(0)h = - \{ \partial_x^2 + r^{-2} \partial_\theta^2 \} \{ \partial_x^2 + r^{-2} \partial_\theta^2 + r^{-2} \} h, \quad h \in bc_{2\pi,per}^{4+\alpha}(C_r).$$

This can be seen as follows, for instance. Since we already know that G is differentiable, the Fréchet derivative $DG(0)$ can be computed as

$$DG(0)h = \left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} G(\varepsilon h) = \left[\left. \frac{d}{d\varepsilon} \right|_{\varepsilon=0} \Delta_{\Gamma(\varepsilon h)} \right] \mathcal{H}(0) + \Delta_{\Gamma(0)} \mathcal{H}'(0)h.$$

Here one notes that $\Delta_{\Gamma(\epsilon h)}\mathcal{H}(0) \equiv 0$ (since $\mathcal{H}(0) = \mathcal{H}_{\mathcal{C}_r} \equiv 1/r$ is constant) while

$$\Delta_{\Gamma(0)} = \Delta_{\mathcal{C}_r} = \partial_x^2 + r^{-2}\partial_\theta^2,$$

and $\mathcal{H}'(0) = -(\Delta_{\mathcal{C}_r} + r^{-2})$ follow from [34, formula (31)] (noting that, by orientation convention, mean curvature has the opposite sign in [34]).

Lemma 4.1. *The spectrum of $DG(0)$ consists entirely of eigenvalues and is given by*

$$(4.2) \quad \sigma(DG(0)) = \{-(m^2 + r^{-2}n^2)(m^2 + r^{-2}n^2 - r^{-2}) : m, n \in \mathbb{Z}\}.$$

Moreover, 0 is a semi-simple eigenvalue of multiplicity 3 and

$$N(DG(0)) = \text{span}\{1, \cos(\theta), \sin(\theta)\}.$$

Proof. Due to compactness of the embedding $bc_{2\pi, \text{per}}^{4+\alpha}(\mathcal{C}_r)$ into $bc_{2\pi, \text{per}}^\alpha(\mathcal{C}_r)$ we know that the spectrum of $DG(0)$ consists only of eigenvalues. Taking advantage of Proposition 3.8, we conclude that

$$\sigma(DG(0)) = \{-(m^2 + r^{-2}n^2)(m^2 + r^{-2}n^2 - r^{-2}) : m, n \in \mathbb{Z}\}.$$

Recalling the expression (4.1) for the linearized evolution operator, $DG(0)$ is realized as a Fourier multiplier acting on $h \in bc_{2\pi, \text{per}}^{4+\alpha}(\mathcal{C}_r)$ as

$$DG(0)h = \sum_{m, n \in \mathbb{Z}} \hat{h}(m, n)[(m^2 + r^{-2}n^2)(m^2 + r^{-2}n^2 - r^{-2})]e^{imx}e^{in\theta}.$$

Locating the kernel of the linearization, if $DG(0)h = 0$, then the fact that $r > 1$ necessarily implies that $\hat{h}(m, n) = 0$ whenever $m \neq 0$, and for $|n| > 1$ when $m = 0$. Thus, the kernel of $DG(0)$ coincides with the 3-dimensional subspace spanned by the first order cylindrical harmonics $\{1, \cos(\theta), \sin(\theta)\}$.

Similarly, if $DG(0)h \in N(DG(0))$, it follows that $\hat{h}(m, n) = 0$ for all $m \neq 0$ and for $|n| > 1$ when $m = 0$. Hence, $N(DG(0)^2) = N(DG(0))$ and 0 is thus a semi-simple eigenvalue of $DG(0)$ with both geometric and algebraic multiplicity 3 and eigenspace N . \square

Note that $DG(0) \subset (-\infty, 0]$ when $r \geq 1$, which we will exploit in order to obtain stability. However, accounting for the presence of 0 in the spectrum $DG(0)$ — and acknowledging the fact that the family of cylinders \mathcal{C}_r are obviously not isolated equilibria — we turn to a center-manifold analysis for the stability analysis.

4.2. Local Family of Cylinders. Before we can apply the main result of [33], we must address the local nature of the family of equilibria surrounding the cylinder. We show in this section that the family of equilibria surrounding $\rho \equiv 0$ is a smooth 3-dimensional manifold in $bc_{2\pi, \text{per}}^{4+\alpha}(\mathcal{C}_r)$. Specifically, we shall show that the local family of equilibria coincides precisely with the manifold of two-dimensional cylinders near \mathcal{C}_r , for which we provide an explicit parametrization.

Denote by $\mathcal{M}_{\text{cyl}} \subset bc^{4+\alpha}(\mathcal{C}_r)$ the family of height functions $\bar{\rho} = \bar{\rho}(\bar{y}, \bar{z}, \bar{r})$ such that

$$\Gamma(\bar{\rho}) = C(\bar{y}, \bar{z}, \bar{r}) := \{(x, y, z) \in \mathbb{R}^3 : (y - \bar{y})^2 + (z - \bar{z})^2 = \bar{r}^2, x \in \mathbb{R}\}$$

i.e. $\Gamma(\bar{\rho})$ is an infinite cylinder with axis of symmetry (x, \bar{y}, \bar{z}) and $\bar{r} > 0$, (in contrast to \mathcal{C}_r which is centered about the axis $(x, 0, 0)$ with radius r). It follows

that $\bar{\rho} = \bar{\rho}(x, \theta)$ must be independent of x , and, as depicted in Figure 1, it is clear that $\bar{\rho} \in \mathcal{M}_{cyl}$ must satisfy the equation

$$((r + \bar{\rho}) \cos(\theta) - \bar{y})^2 + ((r + \bar{\rho}) \sin(\theta) - \bar{z})^2 = \bar{r}^2.$$

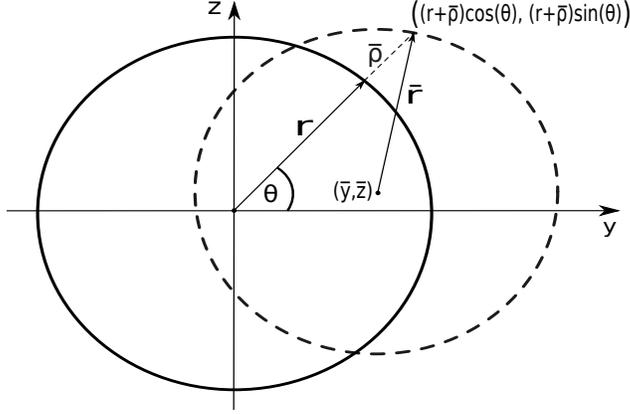


FIGURE 1. Cross section of $C(\bar{y}, \bar{z}, \bar{r})$ defined over \mathcal{C}_r via the height function $\bar{\rho} = \bar{\rho}(x, \theta)$.

Solving this quadratic equation for $(r + \bar{\rho})$ readily yields the formula

$$\bar{\rho}(x, \theta) = \cos(\theta)\bar{y} + \sin(\theta)\bar{z} + \sqrt{\bar{r}^2 - (\sin(\theta)\bar{y} - \cos(\theta)\bar{z})^2} - r.$$

The mapping $[(\bar{y}, \bar{z}, \bar{r}) \mapsto \bar{\rho}]$ is a local diffeomorphism from some neighborhood $\mathcal{O} \subset \mathbb{R}^3$ of $(0, 0, r)$ into $bc_{2\pi, per}^{4+\alpha}(\mathcal{C}_r)$, so that \mathcal{M}_{cyl} is a 3-dimensional Banach manifold of equilibria containing $\rho \equiv 0$.

Next, we wish show that \mathcal{M}_{cyl} contains **all** equilibria of (2.5) in a neighborhood of $\rho \equiv 0$. We accomplish this by showing that \mathcal{M}_{cyl} coincides locally with a stable center manifold for (2.5).

Theorem 4.2. *Given $r > 1$, (2.5) admits a unique, 3-dimensional stable center manifold in a neighborhood of $\rho \equiv 0$. Moreover, this center manifold must coincide with \mathcal{M}_{cyl} , which thus contains all equilibria near $\rho \equiv 0$.*

Proof. For $N = \text{span}\{1, \cos(\theta), \sin(\theta)\}$ let $\pi^c : bc^\sigma(\mathcal{C}_r) \rightarrow N$ be the orthogonal projection, given by

$$\pi^c h := (h|1)1 + (h|\cos(\theta))\cos(\theta) + (h|\sin(\theta))\sin(\theta),$$

where the inner product $(\cdot|\cdot)$ is defined for 2π -periodic, L_2 integrable functions on \mathcal{C}_r . In particular,

$$(h|g) := \int_{\mathbb{T}^2} h(x, \theta)g(x, \theta)dx d\theta$$

compares functions only over one interval of periodicity along the x -axis. The space of periodic L_2 functions over \mathcal{C}_r clearly contains $bc_{2\pi, per}^\sigma(\mathcal{C}_r)$ as a subspace, for any $\sigma > 0$. Defining $\pi^s := id - \pi^c$ and $S := \pi^s(bc_{2\pi, per}^{4+\alpha}(\mathcal{C}_r))$ we get the direct topological decomposition $bc_{2\pi, per}^{4+\alpha}(\mathcal{C}_r) = N \oplus S$ which reduces the operator $DG(0)$, as follows from the fact that π^c commutes with $DG(0)$. We conclude that

π^c coincides with the spectral projection of $DG(0)$ corresponding to the spectral set $\{0\}$. The existence of a (local) center manifold now follows from [37, Theorem 4.1]. In particular, there exists a neighborhood \mathcal{U} of $0 \in N$ and a mapping

$$\psi \in C^k(\mathcal{U}, bc_{2\pi, per}^{4+\alpha}(\mathcal{C}_r)), \quad \text{for any } k \in \mathbb{N} \setminus \{0\},$$

such that $\psi(0) = 0$, $\partial\psi(0) = 0$ and the graph

$$\mathcal{M}^c := \text{graph}(\psi) \subset N \oplus S = bc_{2\pi, per}^{4+\alpha}(\mathcal{C}_r)$$

is a locally invariant manifold for the semiflow generated by (2.5), containing all global small solutions (and, in particular, all equilibria close to $\rho = 0$). Moreover, it follows that the tangent space $T_0(\mathcal{M}^c)$ coincides with the eigenspace N , so \mathcal{M}^c is a 3-dimensional manifold containing $\rho = 0$.

As \mathcal{M}^c contains all equilibria that are close to 0, we know that $\mathcal{M}_{cyl} \subset \mathcal{M}^c$. Since \mathcal{M}_{cyl} and \mathcal{M}^c both are (smooth) manifolds of the same dimension, a standard inverse function theorem argument yields $\mathcal{M}_{cyl} = \mathcal{M}^c$ in a sufficiently small neighborhood of 0. But this implies that every equilibrium close to $\rho = 0$ lies on \mathcal{M}_{cyl} , and thus is a cylinder. \square

4.3. Stability. We are now prepared to prove the main result of this section regarding stability and instability of the family of cylinders \mathcal{C}_r under sufficiently small admissible perturbations exhibiting 2π -axial-periodicity.

Theorem 4.3 (stability and instability of cylinders). *Let $\alpha \in (0, 1)$ and take $\mu \in [1/2, 1]$ so that $4\mu + \alpha \notin \mathbb{Z}$.*

- a) (*Exponential stability*) *Let $r > 1$ and $\omega \in (0, 1 - r^{-2})$ be fixed. Then there exist positive constants $\delta = \delta(\omega)$ and $M = M(\omega)$ such that, given an admissible periodic perturbation*

$$\rho_0 \in V_{\mu, per} := bc_{2\pi, per}^{4\mu+\alpha}(\mathcal{C}_r) \cap [\rho > -r],$$

with $\|\rho_0\|_{bc^{4\mu+\alpha}} < \delta$, the solution $\rho(\cdot, \rho_0)$ exists globally in time, and there exists a $\bar{\rho} \in \mathcal{M}_{cyl}$ such that

$$\|\rho(t, \rho_0) - \bar{\rho}\|_{bc^{4\mu+\alpha}} \leq M e^{-\omega t} \|\rho_0\|_{bc^{4\mu+\alpha}}.$$

Moreover, the radius \bar{r} of the limiting cylinder $\bar{\rho}$ is uniquely determined by preservation of periodic volume; in particular, \bar{r} is determined by the equality $\text{Vol}(\rho_0) = \text{Vol}(\bar{\rho})$.

- b) (*instability*) *Let $0 < r < 1$, then $\rho = 0$ is unstable in the topology of $bc_{2\pi, per}^{4\mu+\alpha}(\mathcal{C}_r)$.*

Proof. (a) Lemma 4.1 and Theorem 4.2 show that \mathcal{C}_r is *normally stable*, i.e. we have with $A = DG(0)$

- (i) near $\rho = 0$ the set \mathcal{E} of equilibria constitutes a C^1 -manifold in $bc_{2\pi, per}^{4+\alpha}(\mathcal{C}_r)$ of dimension 3,
- (ii) the tangent space for \mathcal{E} at 0 is given by $N(A)$,
- (iii) 0 is a semi-simple eigenvalue of A , i.e. $N(A) \oplus R(A) = bc_{2\pi, per}^\alpha(\mathcal{C}_r)$,
- (iv) $\sigma(A) \setminus \{0\} \subset \mathbb{C}_- = \{z \in \mathbb{C} : \text{Re } z < -\omega\}$ for some $\omega > 0$

and exponential stability follows from [33, Theorem 3.1]. To determine the radius \bar{r} of the limit $\bar{\rho}$, note that Theorem 3.7 implies the equation $\text{Vol}(\bar{\rho}) = \text{Vol}(\rho_0)$ must be satisfied. However, there is a one-to-one correspondence between the enclosed volume $\text{Vol}(\bar{\rho})$ (as defined by (3.5)) and the radius \bar{r} of $\bar{\rho}$, hence \bar{r} is uniquely

determined by volume preservation.

(b) Nonlinear instability of cylinders \mathcal{C}_r with radius $0 < r < 1$ has already been shown in [27, Theorem 5.1]. To be precise, the authors proved this result in [27] within the class of *axisymmetric* perturbations of the cylinder. However, these perturbations also reside in the class of axially-definable surfaces, so nonlinear instability of the cylinder has already been established in the current setting. \square

Remarks 4.4. (a) We remark that the proof of the exponential stability result can also be based on a center manifold analysis, as in [18, Theorem 1.2].

(b) There is a striking difference between the situation in [18] and the current situation concerning equilibria for the surface diffusion flow. While it can readily be inferred that (all embedded, compact, closed) equilibria of (1.1) are spheres (single spheres, or multiple spheres of the same radius), the existence and classification of equilibria in the class of axially-definable surfaces is more difficult and complex. In the special case of *axisymmetric and periodic* surfaces, equilibria are characterized by the Delaunay surfaces, see [27]. The general case for *non-axisymmetric* surfaces seems to be wide open.

(c) Theorem 4.3 is interesting on its own, as it characterizes all *axially-definable periodic* equilibria that are close to \mathcal{C}_r for $r > 1$.

(d) We note that the radius of the cylinder $C(\bar{y}, \bar{z}, \bar{r})$ to which perturbations converge (when $r > 1$) is easily determined by the fact that surface diffusion flow preserves enclosed volume. However, we cannot predict the location of the central axis (\bar{y}, \bar{z}) .

5. NEUMANN-TYPE BOUNDARY CONDITIONS

It is also common to consider surfaces satisfying Neumann-type boundary conditions in the axial direction, c.f. [7, 23]. We show in this section that the well-posedness and stability results stated above continue to hold in subspaces of $bc^\sigma(\mathcal{C}_r)$ satisfying such boundary conditions.

Consider the restricted cylinder $\mathcal{C}_{r,a}$ as in (3.4) and define $bc_N^\sigma(\mathcal{C}_{r,a})$ (for $\sigma > 3$) to be the collection of height functions $\tilde{\rho} : \mathcal{C}_{r,a} \rightarrow \mathbb{R}$ satisfying the boundary conditions

$$\partial_x \tilde{\rho}(0+, \theta) = \partial_x \tilde{\rho}(a-, \theta) = 0 \quad \text{and} \quad \partial_x^3 \tilde{\rho}(0+, \theta) = \partial_x^3 \tilde{\rho}(a-, \theta) = 0,$$

for all $\theta \in \mathbb{T}$, and such that $\tilde{\rho}$ is bc^σ regular up to the closed boundaries of $\mathcal{C}_{r,a}$. It follows directly from [2] and [3] that $bc^\sigma(\mathcal{C}_{r,a})$ is a Banach space of functions over the manifold with boundary $\mathcal{C}_{r,a}$. It then follows that $bc_N^\sigma(\mathcal{C}_{r,a})$ is a closed subspace of $bc^\sigma(\mathcal{C}_{r,a})$, and is hence a Banach space.

Before stating results for (2.5) in the setting of Neumann boundary conditions, we first note that $bc_N^\sigma(\mathcal{C}_{r,a})$ is linearly isomorphic to a closed subspace of $bc_{2a,per}^\sigma(\mathcal{C}_r)$ via the even extension operator

$$\rho(x, \theta) = \psi(\tilde{\rho})(x, \theta) := \begin{cases} \tilde{\rho}(2ka - x, \theta) & \text{if } x \in [(2k-1)a, 2ka], k \in \mathbb{Z} \\ \tilde{\rho}(x - 2ka, \theta) & \text{if } x \in [2ka, (2k+1)a], k \in \mathbb{Z}. \end{cases}$$

Thus, all results derived above in the axial periodic setting can also be applied to surfaces satisfying Neumann boundary conditions via restriction.

Proposition 5.1. (*well-posedness for Neumann boundary conditions*) Fix $\alpha \in (0, 1)$, $a > 0$ and take $\mu \in [3/4, 1]$ so that $4\mu + \alpha \notin \mathbb{Z}$. For each initial value

$$\tilde{\rho}_0 \in V_{\mu, N} := bc_N^{4\mu+\alpha}(\mathcal{C}_{r,a}) \cap [\rho > -r],$$

there exists a unique maximal solution

$$\tilde{\rho}(\cdot, \tilde{\rho}_0) \in C_{1-\mu}^1(J(\tilde{\rho}_0), bc_N^\alpha(\mathcal{C}_{r,a})) \cap C_{1-\mu}(J(\tilde{\rho}_0), bc_N^{4+\alpha}(\mathcal{C}_{r,a})).$$

to (2.5). It follows that

$$\mathcal{D} := \bigcup_{\tilde{\rho}_0 \in V_{\mu, N}} J(\tilde{\rho}_0) \times \{\tilde{\rho}_0\}$$

is open in $\mathbb{R}_+ \times V_{\mu, N}$ and the map $[(t, \tilde{\rho}_0) \mapsto \tilde{\rho}(t, \tilde{\rho}_0)]$ defines an analytic semiflow on $V_{\mu, N}$. Moreover, if there exists $0 < M < \infty$ so that, for all $t \in J(\tilde{\rho}_0)$,

- a) $\tilde{\rho}(t, \tilde{\rho}_0)(p) \geq 1/M - r$ for all $p \in \mathcal{C}_{r,a}$, and
- b) $\|\tilde{\rho}(t, \tilde{\rho}_0)\|_{bc_N^{4\mu+\alpha}(\mathcal{C}_{r,a})} \leq M$,

then it must hold that $t^+(\tilde{\rho}_0) = \infty$ and $\tilde{\rho}(\cdot, \tilde{\rho}_0)$ is a global solution.

Proof. Given $\tilde{\rho}_0 \in V_{\mu, N}$, the even extension $\rho_0 := \psi(\tilde{\rho}_0)$ is in the admissible class $V_{\mu, per}$ of $2a$ -periodic functions on \mathcal{C}_r , and hence Proposition 3.6 gives existence, uniqueness, and global well-posedness of a periodic solution

$$\rho(\cdot, \rho_0) \in C_{1-\mu}^1(J(\rho_0), bc_{2a, per}^\alpha(\mathcal{C}_r)) \cap C_{1-\mu}(J(\rho_0), bc_{2a, per}^{4+\alpha}(\mathcal{C}_r))$$

to (2.5) on \mathcal{C}_r . It remains to show that, upon restriction to $\mathcal{C}_{r,a}$, the solution $\rho(t, \rho_0)$ satisfies desired Neumann boundary conditions for all $t \in J(\rho_0)$.

Define the reflection operator R about the $x = 0$ plane so that

$$[R\rho](x, \theta) := \rho(-x, \theta).$$

Then the restriction $\tilde{\rho}(t) := \rho(t)|_{\mathcal{C}_{r,a}}$ satisfies the desired boundary conditions if and only if

$$R\rho_x = \rho_x, \quad R\rho_{xxx} = \rho_{xxx}, \quad R(T_{-a}\rho_x) = T_{-a}\rho_x, \quad \text{and} \quad R(T_{-a}\rho_{xxx}) = T_{-a}\rho_{xxx},$$

where T_{-a} is the axial shift operator discussed in section 3.4.

From a geometric perspective, it is immediately clear that reflections about $x = 0$ commute with the geometric calculations used to compute the evolution operator G , however, confirming commutativity of R and G algebraically requires a little finesse. In particular, note that R commutes with the operations:

$$[\rho \mapsto 1/\rho], \quad [\rho \mapsto \partial_\theta^k \rho], \quad k \in \mathbb{N} \quad \text{and} \quad [(\rho, h) \mapsto \rho h],$$

while

$$R \partial_x^k = (-1)^k \partial_x^k R.$$

However, one should observe that all odd axial derivatives of ρ produced in the operator G are found in products with an even number of similar terms. For instance, ρ_x^2 , $\rho_x \rho_{xxx}$, and $\rho_x \rho_{x\theta}$ are three such products found in the expansion of $G(\rho)$. In this way, one confirms that the operators R and G indeed commute.

Finally, let $\rho_R(t) := R\rho(t)$ and $\rho_{R,a}(t) := RT_{-a}\rho(t)$ for $t \in J(\rho_0)$. Then

$$\begin{aligned} \partial_t \rho_R(t) &= R\partial_t \rho(t) = RG(\rho(t)) = G(R\rho(t)), \quad t \in J(\rho_0), \\ \rho_R(0) &= R\rho(0) = R\rho_0 = \rho_0, \end{aligned}$$

showing that $\rho(\cdot)$ and $\rho_R(\cdot)$ are both solutions of (2.5) with the same initial value. By uniqueness, $R\rho(t) = \rho(t)$, implying that $\rho(t)$ satisfies Neumann boundary conditions at $x = 0$ for all $t \in J(\rho_0)$. Similarly, note that $\rho_{R,a}(t)$ and $T_{-a}\rho(t)$ are both solutions to

$$\partial_t \rho(t) = G(\rho(t)), \quad t \in J(\rho_0), \quad \rho(0) = T_{-a}\rho_0,$$

so that $\rho(t)$ also satisfies Neumann boundary conditions at $x = a$ and thus, by restriction, we have a solution

$$\tilde{\rho}(t) := \rho(t)|_{\mathcal{C}_{r,a}}$$

satisfying Neumann boundary conditions on $\mathcal{C}_{r,a}$, as desired. \square

Remarks 5.2. (a) Regarding statements of stability in the setting of Neumann boundary conditions, we note that stability of cylinders was derived for 2π -periodic perturbations of the unbounded cylinder \mathcal{C}_r . Thus, by fixing $a = \pi$, one can derive stability and instability of the bounded cylinder $\mathcal{C}_{r,\pi}$ under perturbations satisfying Neumann boundary conditions, with stability holding for $r > 1$ and instability for $0 < r < 1$ as before.

(b) We also note that the fundamental interval of periodicity has been taken to be fixed at 2π throughout the article, however this is only a convenience. As discussed in [26, Remarks 6.2], one can take perturbations of any periodicity $a > 0$, and derive analogous stability of cylinders with radius $r > a/2\pi$ and instability for cylinders of radius $0 < r < a/2\pi$.

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